

BiF_3



$V_i (YX_3, \text{zyl } Y = N, P, As, Pb, Sb, Bi, \dots)$
 $X = H, F, Cl, CH_3, C \equiv CH$

Smit W. M. A., and J. H. van
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XIII 936

лсб огул

PX.1968

20

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1971

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Intramolecular force fields in group
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5

$PbCl_3$; $PbBr_3$; PbI_3 ; AsF_3 , $AsCl_3$, $AsBr_3$
 SbF_3 , $AsCl_3$; $AsBr_3$, AsI_3 ; SbF_3 ;
 $BiCl_3$; $SbBr_3$; SbI_3 ; BiF_3 , $BiCl_3$
 $BiBr_3$, BiI_3 (м.н.) XIII 541

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31, №6, 1027-29

PX72

20

(9)

$F_3; AsF_3; SbF_3; BiF_3; BF_3;$
 $PCl_3; AsCl_3; SbCl_3; BiCl_3; BCl_3;$
 $PBr_3; AsBr_3; SbBr_3; BiBr_3; BBr_3;$
 $PI_3; AsI_3; SbI_3; BiI_3; BI_3$

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Тимошин В. С., Тогнев И. Н.,
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Периодические центры асимметрии
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 СД, 1973, 13, № 16, 1022-28 к

Ю (Ф)

Bi F₃

1974

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иттрия и висмута по дан-
ным дифракции нейтронов.

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(см V F₃; T)

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Зак. 247

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1983

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Novak I., Potts A.W.,

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УФ фото-
электро-
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BiF_3

1989

111: 242939r Laser Raman, infrared spectra of bismuth fluoride and its vibrational analysis. Mohan, S.; Kuttiappan, P. (Raman Sch. Phys., Pondicherry Univ., Pondicherry, 605 006 India). *Indian J. Phys., B* 1989, 63B(3), 386-90 (Eng). The IR and polarized Raman spectra of BiF_3 were studied at 0-200 cm^{-1} and vibrational assignments were made along with a normal coordinate anal. using a general quadratic function. Force consts., vibrational amplitudes at 298.16 K, Coriolis coupling consts., centrifugal distortion consts. and potential energy distributions are reported.

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Di

C.A. 1989, 111, N 26

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OM. 36909

1992

геометр.,
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ИК спектры,
св. спектры,
ab initio
расчет

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14 Б1072. Три закона в применении к $D_0(\text{BiF})$. Three laws for $D_0(\text{BiF})$ / Yoo R. K., Ruscic B., Berkowitz J. // Chem. Phys. — 1992. — 166, № 1—2. — С. 215—227. — Англ.

С помощью МС с фотоионизацией и УФ-диапазоне определены кривые выхода фотоионов BiF и BiF_2 , образующихся при испарении смеси $\text{Bi}-\text{BiF}_3$. Найдено, что адiabатич. ПТ ионизации равны $8,658 \pm 0,012$ эВ для BiF и $8,05 \pm 0,05$ эВ для BiF_2 . Порог образования Bi^+ из BiF равен $11,12 \pm 0,05$ эВ; отсюда следует, что энергия диссоциации $D_0(\text{BiF}) \leq 3,84 \pm 0,05$ эВ. Применение второго и третьего законов термодинамики к равновесной р-ции $2\text{Bi}(г) + \text{BiF}_3(г) \rightarrow 2\text{BiF}(г)$ привело к значению $D_0(\text{BiF}) = 3,76 \pm 0,13$ эВ. Аналогично изучение р-ции $\text{Bi}(г) + \text{BiF}_2(г) \rightarrow 2\text{BiF}(г)$ дало след. результаты: $D_0(\text{FBi}-\text{F}) = 3,50 \pm 0,15$ эВ; $D_0(\text{F}_2\text{Bi}-\text{F}) = 4,5 \pm 0,2$ эВ. Обсуждены причины расхождений между полученными в настоящей работе данными и ранее опубликованными результатами. Библ. 47. А. А. Кирюшкин

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1994

BiF_3 u sp.
r. g. u. f. v.
S

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Trends in Inversion Barriers of Group 15 Compounds. 2. Ab-Initio and Density Functional Calculations on Group 15 Fluorides MF_3 (M = N, P, As, Sb, Bi).

BiF3

(Dm: 39074)

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F: BIF3

P: 3

1999

132:16588 Symmetry-broken inversion structures for Group 15 EX3 halides. Schwerdtfeger, Peter; Hunt, Patricia Department of Chemistry, The University of Auckland Auckland, N. Z. Adv. Mol. Struct. Res., 5, 223-262 (English) 1999 A review with 79 refs. The inversion process of EX3 Group 15 hydrides and halides is reviewed. All Group 15 hydrides EH3 and N halides NX3 inv through the classical D3h trigonal planar transition state. For the heav halides, however, the a2 HOMO can interchange with the a1 LUMO as pointed in 1980 by Marynick. Consequently, a low lying 1E' excited state can cou with the 1A1' ground state to undergo an e' 2nd-order Jahn-Teller (SOJT) distortion to a lower lying C2v inversion transition state. Hence, for t heavier Group 15 halides the Dixon-Arduengo

C. A. 2000, 132

edge inversion process through T-shaped transition state is preferred. The potential energy surface (PE for these compds. is analyzed. For PF_3 , AsF_3 , SbF_3 , and BiF_3 at the Hart Fock level the symmetry breaking occurs immediately upon distortion from C_{3v} min. towards the D_{3h} point of the PES. Different topologies around t sym. D_{3h} point are possible in the planar EX_3 arrangement, which connect T- and Y-shaped C_{2v} structures. If e' -SOJT symmetry breaking occurs, the topologies can be derived from either a Mexican hat or a monkey saddle. the high-energy D_{3h} point often cannot be described in a satisfactorily w by single-ref. methods. CASSCF calcns. show significant mixts. between 2 configurations of $1A_1'$ symmetry, 1 which is described by $a_{22}a_1'0$ and deno as $D_{3h}(a_2)$ and the 2nd configuration by $a_{20}a_1'2$ and denoted as $D_{3h}(a_1')$. Electron correlation effects are therefore important for the accurate det of the inversion barrier. For mols. with large configuration mixing (PCl PBr_3 , and PI_3) a definite decision on the symmetry of the inversion struc cannot be made without higher level calcns. For the Bi halides the e' -SO distortion is small despite the energetically more favored $D_{3h}(a_1')$ configuration. For these mols. the PES connecting the planar structures very shallow and a definite decision on the symmetry of the inversion transition state cannot be made. The question whether or not a mol. deco before it inverts is addressed. Common models for the inversion mechanis and bonding for the EX_3 compds. are critically analyzed. Structures, vibrational

F: BiF3

P: 3

134:372013 **Density Functional Studies on the Lone Pair Effect of the Trivalent Group (V) Elements: I. Electronic Structure, Vibronic Coupling, and Chemical Criteria for the Occurrence of Lone Pair Distortions in AX₃ Molecules (A=N to Bi; X=H, and F to I).** Atanasov, M.; Reinen, D. Fachbereich Chemie, Philipps-Universität und Zentrum fuer Materialwissenschaften, Marburg, Germany. J. Phys. Chem. A (2001), 105(22), 5450-5467. in English.

The energetic, steric, and bonding properties of mols. AX₃ (A = N to Bi; X = H, F to I) are analyzed using d. functional theory. It is found that the "lone pair" in the initial D_{3h} geometry is of central atom pz character for the NX₃ and AH₃ mols., whereas it possesses s symmetry in all other cases - here generally with a strong delocalization toward the ligands. The stabilization of the distorted C_{3v} geometry is due mainly to covalency effects, whereas steric interaction forces according to the Gillespie-Nyholm model do not seem to play a significant role. The application of the conventional vibronic pseudo Jahn-Teller coupling

2001



approach (PJT), here for the $D_{3h} \rightarrow C_{3v}$ transition [$A_1' \otimes (\alpha_2'' + \alpha_1') \otimes A_2''$ interaction], is an appropriate means for inorg. chemists to predict trends for the extent of distortion and for the corresponding energy gain. The vibronic coupling const. and the vibronic stabilization energies, which mainly det. the total $D_{3h} \rightarrow C_{3v}$ energy gain, vary according to the sequences $F > H > Cl > Br > I$ (A: N to Bi), and $N > P > As > Sb > Bi$ (X: H,F), the dependence on A being only small or not present (X: Cl to I). Thus, the hardest mols. are the most susceptible to vibronic coupling, the latter energy being approx. imaged by the hardness difference $\eta(C_{3v}) - \eta(D_{3h})$. A roughly inverse trend is obsd. if the extent of the angular distortion $\tau\alpha$ from D_{3h} to C_{3v} symmetry is considered; here, the softest mols. such as $Sb(Bi)Br_3$ exhibit the largest and NH_3 the smallest deviations from D_{3h} geometry.

The different sequences for $\tau\alpha$ are due to the strong influence of the force const., which represents the $C_{3v} \rightarrow D_{3h}$ restoring energy. It is remarkable that the vibronic coupling energy is strongly correlated with the chem. hardness η (an observable quantity), while the stabilization energy for the $D_{3h} \rightarrow C_{3v}$ transition is not directly reflected by η , in contrast to what is generally called the "principle of max. hardness".